

# Snowmass2021 - Letter of Interest

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2 **CF Topical Groups:** (check all that apply /)

3  (CF1) Dark Matter: Particle Like

4  (CF2) Dark Matter: Wavelike

5  (CF3) Dark Matter: Cosmic Probes

6  (CF4) Dark Energy and Cosmic Acceleration: The Modern Universe

7  (CF5) Dark Energy and Cosmic Acceleration: Cosmic Dawn and Before

8  (CF6) Dark Energy and Cosmic Acceleration: Complementarity of Probes and New Facilities

9  (CF7) Cosmic Probes of Fundamental Physics

10  (Other)

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## 13 Silicon Photomultipliers as a target for low-mass dark matter searches

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24 The ability to produce high radiopurity Silicon makes it a desirable candidate for low-  
25 background, rare-event searches. Production R&D on Silicon Photo Multipliers (SiPMs), to  
26 be used in DarkSide-20k, has also yielded new techniques for reducing the inherent noises  
27 associated with SiPMs operating at cryogenic temperatures. We propose a novel SiPM-  
28 based, Silicon target, detector design capable of achieving a very low threshold, on the order  
29 of 150 eV. This design would be sensitive to WIMP-like particles with masses  $< 1$  GeV. A  
30 tower consisting of multiple wafers, along with the associated integrated electronics, would  
31 be constructed using a high-purity copper structure. Housed in a Liquid Neon (LNe) active  
32 veto, allowing the SiPM dark rate to be kept below  $10^{-3}$  cps/mm<sup>2</sup>, a 1 kg Silicon target  
33 could reach a sensitivity better than  $10^{-42}$  cm<sup>2</sup> for both low mass WIMP-like particle and  
34 electrophilic dark matter.

## 1. Introduction

36 Silicon has proven to be an optimal target for dark matter searches at low recoil energy and  
 37 correspondingly low WIMP-like particle mass. The combination of low ionization yield and the  
 38 intrinsically high radiopurity allows it to reach unprecedented background levels [1, 2]. Recent  
 39 results from the Global Argon Dark Matter Collaboration (GADMC) has shown that it is possible to  
 40 develop large area ( $25 \text{ cm}^2$ ) SiPM-based photo-detectors operating in liquid argon with remarkably  
 41 low noises, dark count rate lower than  $10^{-3} \text{ cps/mm}^2$ , direct cross talk (DiCT) probability lower  
 42 than 20 %, and with a signal to noise ratio (SNR) in excess of 20 [3–5].

43 With the knowledge and experience gained through the SiPM R&D process for DarkSide-20k a  
 44 dedicated low-mass WIMP-like particle detector design is presented here. The proposed detector  
 45 would exploit a strong pixelization and an active Liquid Neon (LNe) veto to tag the intrinsic and  
 46 the external background. This detector will be equally sensitive to low mass WIMP-like particle  
 47 and to electrophilic dark matter. Below we provide a sketch of the proposed detector along with  
 48 some very preliminary sensitivity projections.

## 2. Back Side Illuminated SiPMs as an active target for dark matter searches

49 In standard SiPMs, the incoming photons are converted into charge carriers in the shallow depletion  
 50 region that extends for  $5 \text{ }\mu\text{m}$  to  $10 \text{ }\mu\text{m}$  from the top-side of the device. In Back Side Illuminated  
 51 SiPMs (B-SiPM) the back-side is what faces the incoming photons, with the depletion region  
 52 extending for almost the full thickness of the silicon substrate. Typical B-SiPMs have a thickness  
 53 of the order of  $10 \text{ }\mu\text{m}$  to  $50 \text{ }\mu\text{m}$ , however technological processes are available to develop B-SiPMs  
 54 with thicknesses up to  $500 \text{ }\mu\text{m}$  with a cell size of a few micrometers. Assuming that the intrinsic  
 55 SiPMs noises, i.e. DCR and DiCT, are kept to minimum, such B-SiPMs could be an interesting  
 56 target in experimental searches for low-mass WIMP-like particles.

58 In SiPMs operating below 100 K, the rate of dark counts due to the thermal generation of carriers  
 59 is suppressed by more than 7 orders of magnitude with respect to room temperature [3]. Also at this  
 60 temperature, the main contribution to the DCR is generated in the high field region by the carriers’  
 61 tunneling. In B-SiPMs the design of the high field region remains unaltered, so it is reasonable to  
 62 assume that the DCR will not increase. However, special attention will have to be payed in order  
 63 to avoid increasing the cross-talk probability, and this will be the main R&D challenge.

64 A charged particle interacting in the B-SiPM’s bulk will generate about  $250 e^-/\text{keV}$  [6]. In the  
 65 depletion region the recombination probability is low enough to preserve the carriers during their  
 66 drift toward the multiplication region of the cell. Since the resulting signal amplitude is proportional  
 67 to the number of cells reached by the carriers, and not the amount of charge arriving at a particular  
 68 cell, electrons will not disrupt the signals. On average the electrons will drift for  $200 \text{ }\mu\text{m}$  in a field  
 69 of  $\approx 100 \text{ V/cm}$  and, according to transverse diffusion models [7], we can predict that a point like  
 70 interaction of  $100 \text{ eV}$  will affect a minimum of 4 cells ( $3 \text{ }\mu\text{m}$  to  $5 \text{ }\mu\text{m}$ ). The transverse diffusion can  
 71 be further increased by applying a small magnetic field perpendicular to the drift direction. For  
 72 electromagnetic interactions releasing an energy above few keV, the ionization cloud will extend  
 73 beyond the SPAD size, significantly increasing the number of affected cells and the signal amplitude.

## 3. Detector design

74 The core of a detector could consist of a monolithic  $14 \times 14 \text{ cm}^2$  silicon wafer, containing about 280  
 75 B-SiPMs with an area of  $70 \times 1 \text{ mm}^2$  and a thickness of the order of  $500 \text{ }\mu\text{m}$ . The whole unit(die)  
 76 can be directly grown and diced from a single  $8''$  wafer. Based on the results of SiPM production  
 77 for DarkSide-20k we expect the fraction of working pixels on each die to be higher than 85 %.

79 Each B-SiPM die has an active mass of 20 g. In order to increase the detector exposure, several  
 80 dies will be mounted on a high radiopurity copper tower, accruing a total active mass of 1 kg of  
 81 Silicon. The tower will be hosted in a LNe veto, at about 30 K, and at this temperature the SiPMs’  
 82 dark noise rate is expected to be below  $0.5 \text{ mcps/mm}^2$ . A “blackout layer” deposited on the SiPMs’  
 83 surfaces will prevent any undesired detection of photons from Neon scintillation. This protective

84 blackout layer will cover roughly 95 % of the active surface of the detector. The remaining non-blind  
85 pixels would then exploit the Neon scintillation, effectively transforming the LNe passive buffer into  
86 an active veto.

87 The high pixelization of the detector (around 14000 channels) will also be useful in tagging non-  
88 local events and in searching for delayed coincidences, like the  $^{32}\text{Si}-^{32}\text{P}$ .  $^{32}\text{Si}$  is naturally present in  
89 Silicon with a coincidence window of 14 d with a Q-value of 224 keV. Assuming a DiCT probability  
90 of 15 %, the DCR pile-up in  $3\tau$  would be less than a few percent.

91 Integrated electronics can be developed on two  $14\times 5\text{ cm}^2$  monolithic silicon boards to be wire  
92 bonded to the SiPM array. The BNL group has already shown, in their work on DUNE, the  
93 working design of low-power, low-noise, charge amplifiers which feature on-chip digitization and  
94 operate successfully in liquid Argon [8]. And, the lower operating temperature of LNe is not  
95 expected to pose any issues for the CMOS electronics. The boards will be daisy-chained by means  
96 of high radiopurity flat cables, and only a single connection would be required to extract the timing  
97 and charge information for all of the hits in the detector.

98

#### 4. Physics reach

99 Owing to the high radiopurity of the Silicon substrate, the radioactive background in the wafers  
100 is highly, which suppressed. The largest background for this detector comes from the cross-talk,  
101 which enhances the dark rate signal to many photoelectrons. This background follows a Poissonian  
102 distribution which is scaled by the total DCR and fewer than 4 counts/yr would have an ampli-  
103 tude larger than 8 photoelectrons. With the previous assumptions on carriers diffusion, a simple  
104 toy montecarlo on transverse diffusion predicts that a point like interaction of 8 photoelectrons,  
105 corresponding to a visible energy of about 150 eV, would have an acceptance above 80 %. We also  
106 note that this could likely be reduced to less than 100 eV by means of a transverse magnetic field.  
107 Electromagnetic interactions with a deposited energy a few keV will result in events with a large  
108 signal outside the region of interest. The same applies for  $\alpha$ 's decays with typical energy in the  
109 MeV scale. Partially contained electromagnetic or  $\alpha$ 's interactions will be tagged by the active LNe  
110 veto.

111 The detector would require installation in an underground laboratory to reduce the rate of  
112 cosmogenic particles, and a passive shield composed of ultra-pure water and plastic layers would  
113 be required to attenuate the residual neutrons and the gamma rays. The careful estimation of  
114 the background will require a complete simulation of the detector materials, of the shielding, of  
115 the underground laboratory, and of the multi-scatter interactions within the veto and the detector  
116 itself. However, the intrinsic constituents of the detector are very radiopure: silicon, copper and  
117 neon are proven to reach a contamination level of the ppb scale for the Uranium and Thorium  
118 chains. Furthermore, all of the non-stable isotopes of neon have a very short lifetime. Assuming  
119 that the remaining background sources can be kept below the dark rate contribution, we can project  
120 our sensitivity better than  $10^{-42}\text{ cm}^2$  for both low mass WIMP-like particle and electrophilic dark  
121 matter.

122 Furthermore, the anisotropic structure of the semiconductor crystal makes the ionization thresh-  
123 old sensitive to the recoil direction [9]. This produces a modulation in the rate of dark matter  
124 events above the detector energy threshold, which can be as large as  $\pm 40\%$  for MeV-GeV mass  
125 dark matter [10]. A detector configuration where sets of SiPM are arranged in orthogonal planes  
126 enhances the sensitivity to a sidereal modulation signal.

127

#### 5. Conclusion

128 A SiPM-based detector using high a radiopurity Silicon target operating in a LNe active veto could  
129 achieve a threshold of  $\approx 150\text{ eV}$ . This design would be sensitive both to WIMP-like particles and  
130 electrophilic dark matter and allow for further exploration of the parameter space below  $10\text{ GeV}/c^2$ .  
131 The modest production cost of SiPM wafers and the proposed single-cable readout system allow  
132 to easily scale-up the target mass of the detector, making it very competitive with respect to  
133 CCD-based technologies and allowing a cross-check of the current silicon best limits [2].

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