Snowmass2021 - Letter of Interest

² Resonant Halo Axion Detectors for the Mass Range ³ 16-41 µeV

Thematic Areas: (check all that apply \Box/\blacksquare)

□ (CF1) Dark Matter: Particle Like ■ (CF2) Dark Matter: Wavelike

□ (CF3) Dark Matter: Cosmic Probes

- □ (CF4) Dark Energy and Cosmic Acceleration: The Modern Universe
- □ (CF5) Dark Energy and Cosmic Acceleration: Cosmic Dawn and Before
- CF6) Dark Energy and Cosmic Acceleration: Complementarity of Probes and New Facilities
- □ (CF7) Cosmic Probes of Fundamental Physics
- (IF01) Quantum Sensors

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Abstract: We present a summary of efforts to probe the mass range $16 - 41 \,\mu eV$ for dark matter halo axions. Technological challenges in probes of this mass range are discussed, and some solutions outlined. Solutions to some of these challenges are also of interest to communities involved in more general quantum measurement, such as quantum information science. We outline the experimental challenges anticipated when probing this mass range and summarize experimental efforts in which the authors plan to search this mass range for axions in the next 5-10 years.

4 1 Introduction

Axions, representing the potential solution of the 44-year-old strong CP Problem¹, are a candidate of rapidly 5 increasing interest in the quest to solve the 90-year-old dark matter problem^{2,3}. QCD axions⁴ possess 6 very faint couplings to ordinary matter. Search experiments achieving the requisite sensitivity exploit high-7 Q resonators, in which photons from Primakov conversion of axions drive modes of an electromagnetic 8 resonator of volume V threaded by a static magnetic field $B^{5,6}$. When the frequency of a suitable cavity q mode matches the axion energy, resonant enhancement boosts the otherwise undetectable signal power to 10 of order yoctowatts, detectable with existing technology. In practice, the mode frequencies are tuned to 11 probe a range of axion masses at a rate limited by the radiometer equation. Though the range of possible 12 axion masses is many orders of magnitude, recent developments in lattice QCD^{7,8} predict m_a in the range 13 $16 - 41 \,\mu eV$. Some post-inflation production string and domain wall models claim very precise predictions 14 of m_a also within this range⁹. This letter summarizes experimental issues relevant to searches for axions in 15 this well-motivated range of masses. 16

17 2 Magnets and Resonators

Because the signal power is proportional to the magnetic field squared and volume, the first experiments 18 optimized for both B^2V and cost. The ADMX experiment currently uses a single cavity in a 220-liter, 19 7-T solenoidal magnet. Its most recent experiments probe the DFSZ^{10,11} and KSVZ^{12,13} axion models in 20 the range $2.6 - 3.3 \,\mu eV^{14}$. The Haystac haloscope, operating a higher frequencies, uses a smaller volume, 21 9-T solenoidal magnet. Its first results place limits on axions at about twice the KSVZ model power level 22 for a $23.55 - 24.0 \,\mu\text{eV}$ mass range. Probing much higher masses naturally leads to ever-decreasing cavity 23 volumes, and hence less signal power. Several approaches to achieving higher sensitivity in this mass range 24 are possible. First, arrays of small cavities power-combined can be tuned in-phase to achieve the required 25 volume. This difficult technical task has been made tractable in the past decade through the development 26 of cryogenic piezo crystal actuators. Second, a higher field magnet can be used at a smaller volume. De-27 velopments in magnet technologies now enable fields exceeding 30 T over suitable geometries, though such 28 magnets are very costly¹⁵. An aspiration for the future would be funding for a large (order 1m bore) magnet 29 with a high field (order 20-30 T), though such a magnet would itself be a major research and development 30 project. 31

32 **3** Low Noise Receivers

Improvements in low noise amplifier technology may mitigate some of the limitations of axion searches 33 in this mass range. The advent of the Josephson Parametric Amplifier (JPA) has already enabled higher 34 frequency axion searches¹⁴. The HAYSTAC experiment used a JPA in its 23.55–24.0 µeV run¹⁶. The most 35 recent ADMX run also used a JPA to reach DFSZ^{10,11} sensitivity¹⁴. Traveling wave parametric amplifiers 36 (TWPAs) may also prove beneficial in this frequency regime¹⁷. A major advantage of TWPAs is that they 37 do not rely on a resonant structure and therefore provide gain over a wide bandwidth, reducing the need to 38 tune the amplifier. They are also less sensitive to perturbations in the magnetic field. These features could 39 help reduce some of the complexity that is inherent to multi-cavity systems that may be necessary in this 40 frequency range. Such devices can provide a power gain of 20 dB over a bandwidth of 3 GHz^{17,18}. Quantum 41 linear amplifiers such as JPAs and TWPAs can achieve exceedingly low noise levels, but are constrained to 42 operate above the standard quantum limit. HAYSTAC has recently used the enhancement from a squeezed 43

state receiver to double their search rate, setting a new limit in the region from 16.96–17.12 and 17.14– 44 $17.28 \ \mu eV^{19}$. Single photon counters, if feasible in this frequency range, would provide an alternative to 45 quantum linear amplifiers that could radically change the landscape of possible experiments. In fact, there 46 is reason to believe that for a least part of the $16 - 41 \,\mu eV$ mass range, single photon detectors would be 47 preferable to linear amplifiers²⁰. There is some evidence that suggests this may be possible in the near 48 future²¹. Single photon counters using Josephson Junctions have shown some promise²², but further studies 49 are necessary before this technology can be fully implemented. An approach pioneered by the CARRACK 50 collaboration^{23,24} coupled single photon detectors in the form of rubidium Rydberg atoms^{23,24} to an axion 51 haloscope, excluding axions of mass $\sim 1.006 \times 10^{-5} \,\mathrm{eV}$ at the level of the DFSZ model band. The continued 52 investigation of Rydberg atoms as single-photon detectors is being undertaken at Johns Hopkins. 53

54 **4** Experiments and New Target Technologies

Current and near-future ADMX operations will cover the mass range from $4.1 \,\mu eV$ to $16 \,\mu eV$ using multi-55 cavity arrays that fit within the confines of the existing ADMX magnet bore, or other readily available 56 options. Going to higher masses requires either the use of a large number of smaller cavities, or a higher 57 magnetic field, or both. A UK-based group collaborating with ADMX is now in a position to develop a 58 two-pronged approach to cover this frequency range. First, they will collaborate with ADMX to append an 59 additional receiver chain to the existing insert RF electronics. Second, it will also build a new UK-based 60 high-field, high-volume, low-temperature facility for further axion and hidden sector searches in this mass 61 range, in collaboration with ADMX. Another possibility means of scanning this regime is to run a cavity 62 array without attempting to synchronize the resonant frequencies of all the cavities, instead combining the 63 data from all the cavities running at different resonant frequencies to achieve equivalent sensitivity overall. 64 This is the approach adopted in ORGAN²⁵. A similar idea aims to fabricate cavities of varying dimensions 65 coated in superconducting film operated high magnetic fields. Although operating superconducting cavities 66 in high fields is a known technological challenge, there is a strong incentive to develop such technology for 67 Ouantum Information Science (OIS). Another recently proposed alternative approach is to use an external 68 feedback loop incorporating a resonant circuit to induce artificial high Q resonances in an electromagnetic 69 structure²⁶. A potential advantage of this technique is that many resonances could be induced in parallel, 70 resulting in a speed-up in the coverage of mass range by a factor of the number of resonances. Early tests of 71 this approach are underway at ADMX. Above $41 \,\mu eV$ alternative detection schemes such as phased arrays of 72 mirror reflectors as proposed by MADMAX²⁷ or through coupling of higher mass QCD axions to fermionic 73 spin as in QUAX²⁸. However, between 16 μeV and 41 μeV , resonant structures in magnetic fields remain 74 the only viable technology. 75

76 **5** Conclusion

Axions remain one of the most compelling solutions to the dark matter and strong CP problem, with the 16 – 41 µeV mass range well-motivated by current theory. The outstanding experimental issue in this mass range is combining a large volume and 10 - 40 GHz resonant frequencies. We believe that this issue will be the prime focus of the experimental efforts in probing the 16 - 41 µeV mass range. This work will also act as a catalyst for the additional development of quantum electronics at higher frequencies. Further advancements in JPAs and TWPAs will be accompanied by new detectors based on QUBIT arrays. Modern mechanically pre-cooled dilution refrigerators should greatly simplify the cryogenics of new apparatus. We look forward to the ever more sensitive searches, and to the long awaited discovery of axion dark matter.

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