

CP-violation study in *b*-baryon hadronic decays using $SU(3)$ symmetry

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A significant effort has been made by CDF and LHCb experiments towards establishing *CP* violation in hadronic *b*-baryon decays. Future experimental runs, led by the LHCb collaboration, promise a much larger dataset of such decays paving the way for several interesting *CP*-asymmetry measurements. In this letter of interest, we identify some of those decay modes and discuss *CP*-asymmetry relations among them.

I. INTRODUCTION

The *CP*-asymmetry study of bottom baryons [1, 2] is of great interest as it provides a direct probe of matter-antimatter asymmetry in baryon decays. Despite of the experimental challenges, both CDF [3, 4] and LHCb [5–11] have come tantalizingly close to measuring *CP*-violation in *b*-baryon decays. So far, within the current uncertainties, *CP*-asymmetry studies [3, 4, 9] in $\Lambda_b \rightarrow p^+ K^-$ and $\Lambda_b \rightarrow p^+ \pi^-$ decays are consistent with no *CP*-asymmetry. The situation may change, however, when the statistical errors go down as LHCb gears up to analyze a significant number [12] of hadronic *b*-baryons decays in its subsequent runs [13].

One expects *b*-baryon *CP*-asymmetries [14, 15] to be of similar magnitude as observed in *B*-meson decays since the underlying quark level transition are the same for the two cases. Moreover, a non-vanishing *CP*-asymmetry is a measure of direct *CP* violation as baryons and antibaryons do not undergo mixing due to baryon number conservation. The two body [16–20] and multibody charmless decays [6–8, 10, 11, 21, 22] of *b*-baryons mediated by $b \rightarrow u\bar{u}s$ and $b \rightarrow u\bar{u}d$ transitions are ideal places to look for *CP*-violating effects by either measuring decay rate asymmetries directly or through *T*-odd asymmetry observables [23–25] between decay mode and its conjugate. Motivated by the discovery potential of *CP*-violation in baryons, we focus on *b*-baryons decaying to an octet or decuplet baryon and a pseudoscalar meson and use $SU(3)$ -flavor symmetry

to derive *CP*-relations [17, 18, 20] that can be tested in near future [26].

II. $SU(3)$ -DECOMPOSITION OF DECAY AMPLITUDES

The number of $SU(3)$ -reduced matrix elements [18, 20] pertaining to describe a $SU(3)$ -flavor anti-triplet *b*-baryon (\mathcal{B}_b) decaying to a ground state baryon and a pseudoscalar meson (\mathcal{P}) is ten or five depending on whether the daughter baryon belongs to an octet or a decuplet of $SU(3)$. In both these cases, the total numbers of possible strangeness changing ($\Delta S = -1$) and strangeness conserving ($\Delta S = 0$) decay modes are higher than the number of independent $SU(3)$ -reduced matrix elements. As a consequence, several amplitude relations exist between decay modes. *CP* relations can be inferred by noting that both the tree and penguin part of the decay amplitude satisfy the amplitude relations. More than one partial waves can contribute to such decays [1, 2, 27, 28] and the decay amplitude is decomposed as,

$$\mathcal{A}^l = \lambda_u^q \mathcal{A}_{\text{tree}}^l + \lambda_t^q \mathcal{A}_{\text{penguin}}^l. \quad (1)$$

where $l = 0, 1, 2$ denote contributions from particular partial waves. The δ_{CP}^l for each partial wave is given by,

$$\delta_{\text{CP}}^l(\mathcal{B}_b \rightarrow \mathcal{B}\mathcal{P}) = -4\mathbf{J} \times \text{Im} \left[\mathcal{A}_{\text{tree}}^{l*}(\mathcal{B}_b \rightarrow \mathcal{B}\mathcal{P}) \mathcal{A}_{\text{penguin}}^l(\mathcal{B}_b \rightarrow \mathcal{B}\mathcal{P}) \right], \quad (2)$$

\mathbf{J} being the well known Jarlskog invariant and $\mathcal{B} = \mathcal{D}$ or \mathcal{O} depending on whether the final state baryon belongs to a decuplet or an octet of $SU(3)$. The experimentally measured quantity, A_{CP} , is the sum of *CP* violating

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contributions from the δ_{CP}^l with appropriate phase-space factor multiplied. Starting from a Λ_b^0 baryon in the initial state, some of the δ_{CP} relations [20] having decuplet baryon in the final state are quoted below,

$$\begin{aligned} \delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Delta^+ K^-) &= \delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Delta^0 \bar{K}^0) & (3) \\ \delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Sigma'^- \pi^+) &= \delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Xi'^- K^+) \\ &= -\frac{1}{3} \delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Delta^+ \pi^-) = -\delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Sigma'^- K^+) & (4) \end{aligned}$$

In case of octet baryons in the final state, all of the δ_{CP} -relations [18] can be derived from U -spin symmetry,

$$\begin{aligned} \delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Sigma^- K^+) &= -\delta_{\text{CP}}^l(\Xi_b^0 \rightarrow \Xi^- \pi^+), \\ \delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow p^+ \pi^-) &= -\delta_{\text{CP}}^l(\Xi_b^0 \rightarrow \Sigma^+ K^-), \\ \delta_{\text{CP}}^l(\Xi_b^- \rightarrow n K^-) &= -\delta_{\text{CP}}^l(\Xi_b^- \rightarrow \Xi^0 \pi^-), \\ \delta_{\text{CP}}^l(\Xi_b^- \rightarrow \Xi^- K^0) &= -\delta_{\text{CP}}^l(\Xi_b^- \rightarrow \Sigma^- \bar{K}^0), \\ \delta_{\text{CP}}^l(\Xi_b^0 \rightarrow \Xi^- K^+) &= -\delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Sigma^- \pi^+), \\ \delta_{\text{CP}}^l(\Xi_b^0 \rightarrow \Sigma^- \pi^+) &= -\delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Xi^- K^+), \\ \delta_{\text{CP}}^l(\Xi_b^0 \rightarrow \Sigma^+ \pi^-) &= -\delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow p^+ K^-), \\ \delta_{\text{CP}}^l(\Xi_b^0 \rightarrow n \bar{K}^0) &= -\delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Xi^0 K^0), \\ \delta_{\text{CP}}^l(\Xi_b^0 \rightarrow p^+ K^-) &= -\delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow \Sigma^+ \pi^-), \\ \delta_{\text{CP}}^l(\Xi_b^0 \rightarrow \Xi^0 K^0) &= -\delta_{\text{CP}}^l(\Lambda_b^0 \rightarrow n \bar{K}^0), & (5) \end{aligned}$$

Since, A_{CP} depends on the masses of the initial and final baryons as well as the outgoing meson [1, 2], some approximation is needed to obtain A_{CP} relations between various modes. In the U -spin limit, CP violation relations can be experimentally verified using the relation [29],

$$\frac{A_{\text{CP}}(\mathcal{B}_{b,i} \rightarrow \mathcal{B}_j \mathcal{P}_k)}{A_{\text{CP}}(\mathcal{B}_{b,l} \rightarrow \mathcal{B}_m \mathcal{P}_n)} \simeq -\frac{\tau_{\mathcal{B}_{b,i}} \mathcal{BR}(\mathcal{B}_{b,l} \rightarrow \mathcal{B}_m \mathcal{P}_n)}{\tau_{\mathcal{B}_{b,l}} \mathcal{BR}(\mathcal{B}_{b,i} \rightarrow \mathcal{B}_j \mathcal{P}_k)}, \quad (6)$$

where i, j, k and l, m, n are indices corresponding to the various hadrons belonging to the above mentioned δ_{CP} relations. There is a further simplification in case $i = l$, resulting in

$$\frac{A_{\text{CP}}(\mathcal{B}_{b,i} \rightarrow \mathcal{B}_j \mathcal{P}_k)}{A_{\text{CP}}(\mathcal{B}_{b,i} \rightarrow \mathcal{B}_m \mathcal{P}_n)} \simeq -\frac{\mathcal{BR}(\mathcal{B}_{b,i} \rightarrow \mathcal{B}_m \mathcal{P}_n)}{\mathcal{BR}(\mathcal{B}_{b,i} \rightarrow \mathcal{B}_j \mathcal{P}_k)}, \quad (7)$$

where the uncertainties due to lifetime measurement also cancel out [17, 29]. These ratios serve as an important test of flavor $SU(3)$ symmetry in beauty-baryon non-leptonic decays and one can compare these findings with the analogous decays of bottom mesons to have a better understanding of the $SU(3)$ flavor symmetry breaking pattern.

III. FUTURE DIRECTIONS

A non-zero CP asymmetry measurement in b -baryons look promising and it is being pursued alongside B -meson decays at the LHCb. Such observations may provide complimentary information about the CP violating parameter in the quark sector of the Standard Model. While we have focused on two body b -baryon decays, one can also measure CP asymmetry in the final multibody state [6–8, 10, 11] which is often dominated by quasi two-body decays of a b -baryon into another baryon and vector meson. Measurements of CP -violation in regions of phase space using T -odd observables [23–25] have already been performed and new results with updated statistics are awaited.

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